

Spatial Families of Orbits in 2D Conservative Fields

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Abstract. In the framework of the 3D inverse problem of dynamics, we establish the conditions which must be fulfilled by a spatial family of curves to possibly be described by a unit mass particle under the action of a 2D potential $V = v(y, z)$, and give a method to find the potential.

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INTRODUCTION

In the 3D inverse problem of Dynamics, assuming that a two-parametric set of orbits

$$f(x, y, z) = c_1, \quad g(x, y, z) = c_2 \quad (1)$$

in the $Oxyz$ space can be traced by a unit mass particle in the presence of an unknown potential $V = V(x, y, z)$, one aims to finding the potential (see [1] for a review of the results up to 1995).

Bozis and Kotoulas [2] and also Anisiu [3] produced a system of two linear in $V(x, y, z)$ PDEs, one of the first and one of the second order, which will be used in what follows.

We shall treat the problem for the special case of potentials $V = v(y, z)$.

THE EQUATIONS OF THE PROBLEM

We deal with two-parametric families of orbits written in the form (1), which are in an one-to-one correspondence with a pair (α, β) of ‘slope functions’ defined by

$$\alpha = \frac{f_z g_x - f_x g_z}{f_y g_z - f_z g_y}, \quad \beta = \frac{f_x g_y - f_y g_x}{f_y g_z - f_z g_y}. \quad (2)$$

The indices denote partial derivatives.

Let us assume that $\alpha_0 \neq 0$ and adopt the notation

$$\bar{\epsilon} = (1, \alpha, \beta), \alpha_0 = \bar{\epsilon} \text{ grad } \alpha, \beta_0 = \bar{\epsilon} \text{ grad } \beta, \Theta = 1 + \alpha^2 + \beta^2, n = \frac{\Theta}{\alpha_0}, n_0 = \bar{\epsilon} \text{ grad } n. \quad (3)$$

We shall consider exclusively potentials of the form $V = v(y, z)$ and families (1) with $\alpha_0 \neq 0$. In this case the system of the two PDEs mentioned in the introduction becomes

$$v_z = Gv_y, \Theta(\alpha + \beta G)v_{yy} + \Psi v_y = 0, \quad (4)$$

where $G = \beta_0/\alpha_0$, $\Psi = \beta\Theta G_y + \alpha_0(n_0 - 2(\alpha + \beta G))$.

For any compatible pair of potential $v(y, z)$ and orbit (α, β) real motion is allowed in the region $-v_y/\alpha_0 \geq 0$ ([2], [3]).

COMPATIBLE POTENTIALS $V = v(y, z)$ AND FAMILIES (1)

It is seen from (4a) that the function G must be independent of x , i. e.

$$G_x = 0. \quad (5)$$

We assume that $\alpha + \beta G \neq 0$ and we put $H = \Psi/\Theta(\alpha + \beta G)$. For the PDE (4b), now written as $v_{yy} + Hv_y = 0$, to have a solution of the form $v(y, z)$ it must be

$$H_x = 0. \quad (6)$$

From (4b) we get

$$v_y = D(z) \exp\left(-\int^y H(u, z) du\right) \quad (7)$$

and from (4a) we obtain v_z . From the compatibility condition ($v_{yz} = v_{zy}$) for v_y and v_z , as these are given by (7) and (4a) respectively, there follows the homogeneous linear first order ODE $D'(z) = JD(z)$, where J depends merely on z if and only if

$$G_{yy} - G_y H - G H_y + H_z = 0. \quad (8)$$

Proposition For $\alpha_0 \neq 0$, $\alpha + \beta G \neq 0$ and for any family (α, β) satisfying the conditions (5), (6) and (8), there exists a two-dimension compatible potential $v = v(y, z)$. The potential is given by the (compatible) equations (7) and (4a). Real motion is allowed in the region defined by the inequality $D(z)/\alpha_0 \leq 0$.

If $\alpha_0 \neq 0$, $\alpha + \beta G = 0$ and $\Psi \neq 0$ it follows that $v = \text{const}$. If $\alpha_0 \neq 0$, $\alpha + \beta G = 0$ and $\Psi = 0$, equation (4b) is satisfied identically and the potential is found from (4a), provided that the condition (5) holds, and it will not be uniquely determined.

Example For the family $f(x, y, z) = x^4 y z^3$, $g(x, y, z) = x^2 y z$, we get $\alpha_0 \neq 0$, $\alpha + \beta G \neq 0$ and obtain the compatible potential $V(x, y, z) = -(y^2 + z^2)^2$.

The case $\alpha_0 = 0$ and that of potentials depending on (x, y) or (x, z) will be treated elsewhere.

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