

A SIMPLIFIED HOMOTOPY PERTURBATION METHOD
FOR NONLINEAR ILL-POSED OPERATOR EQUATIONS
IN HILBERT SPACES

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Abstract. One popular regularization technique for handling both linear and nonlinear ill-posed problems is Homotopy perturbation. In order to solve nonlinear ill-posed problems, we investigate an iteratively-regularized simplified version of the Homotopy perturbation approach in this study. We examine the method's thorough convergence analysis under typical circumstances, focusing on the nonlinearity and the convergence rate under a Hölder-type source condition. Lastly, numerical simulations are run to confirm the method's effectiveness.

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1. INTRODUCTION

There are numerous scientific and engineering applications that can lead to inverse problems. When inverse issues are formulated mathematically, they typically result in ill-posedness in the sense of Hadamard. Therefore, to find a stable approximation solution for the inverse issue, the regularization method is necessary. Numerous regularization techniques, including the Gauss-Newton iterative approach, the Thikhonov method, the Levrentiv iterative method, and the Levenberg Marquardt method, have been developed in the literature to address issues of this nature in Hilbert spaces (see, for example, [1], [2], [5], [9], [11], [15], [16], and [17]). One of the most popular iterative techniques for resolving nonlinear ill-posed problems in Hilbert spaces is the Landweber technique (cf. [3], [7], [10], [20], and [21]). Compared to other regularization techniques, this one is simpler to implement. A detailed explanation of this technique for the linear case can be found in [4].

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In order to comprehend this approach, let us look at the abstract operator equation

$$(1) \quad F(x) = y,$$

In this equation, $F : D(F) \subseteq X \rightarrow Y$ is a nonlinear operator on $D(F)$, and X and Y are Hilbert spaces with inner product $\langle \cdot, \cdot \rangle$ and norm $\|\cdot\|$ respectively. These operators can always be recognized by their context. $F'(x)$ represents the Fréchet derivative of F at x . Assume for the moment that x^\dagger is the exact solution (which does not need to be unique) to (1.1). We are mainly interested in problems of the form (1) for which the solution x^\dagger does not depend continuously on the data y . In real-world applications, precise data might not be accessible. Therefore, instead of using the actual data y , we use the available perturbed data with

$$(2) \quad \|y^\delta - y\| \leq \delta,$$

where $\delta > 0$ is the noise level.

Given that F is locally scaled and Fréchet differentiable, let

$$(3) \quad \|F'(x)\| \leq 1, \quad \forall x \in B_\rho(x_0),$$

where x_0 is the initial guess for the exact solution and $\rho > 0$. Next, Hanke *et al.* [7] examined the standard Landweber iteration for the non-linear case, which is as follows:

$$(4) \quad x_{n+1}^\delta = x_n^\delta - F'(x_n^\delta)^*(F(x_n^\delta) - y^\delta),$$

where $F'(x)^*$ indicates the adjoint of the Fréchet derivative $F'(x)$ for $x \in B_\rho(x_0)$, and x_0 is the initial guess for the exact solution. To examine technique (4), they took into account the following tangential type nonlinearity condition: With $\eta < \frac{1}{2}$ and $x, x' \in B_\rho(x_0) \subset D(F)$,

$$(5) \quad \|F(x) - F(x') - F'(x)(x - x')\| \leq \eta \|F(x) - F(x')\|.$$

The method's convergence analysis was conducted with the condition (5). For $0 \leq n \leq n_\delta$, the iteration is terminated at n_δ by the generalised discrepancy principle

$$(6) \quad \|F(x_{n_\delta}^\delta) - y^\delta\| \leq \tau \delta < \|F(x_n^\delta) - y^\delta\|, \text{ for } 0 \leq n \leq n_\delta,$$

where $\tau > 2\frac{1+\eta}{1-2\eta} > 2$ is the positive constant dependent on η . Under the subsequent Hölder-type source condition, they were able to determine the rate of convergence:

$$(7) \quad x_0 - x^\dagger = (F'(x^\dagger)^* F'(x^\dagger))^\nu w, \quad w \in X, \quad 0 < \nu \leq \frac{1}{2}.$$

To determine the rate of convergence for the approach (4), they also needed the following features of F , as the assumption (7) is insufficient:

$$(8) \quad F'(x) = R_x^{x'} F'(x'), \quad x, x' \in B_\rho(x_0),$$

where $\{R_x^{x'} : x, x' \in B_\rho(x_0)\}$ is the family of bounded linear operators $R_x^{x'} : Y \rightarrow Y$ such that

$$(9) \quad \|R_x^{x'} - I\| \leq C\|x - x'\|,$$

where C is a positive constant.

Li Cao, Bo Han, and Wei Wang initially created homotopy perturbation iteration for nonlinear ill-posed problems in Hilbert spaces (cf. [13], [14]). The main idea behind it is to integrate the homotopy methodology with the standard perturbation method, and incorporate an embedding homotopy parameter. With $T_n = F'(x_n^\delta)$ as the notation The formula for the N -order homotopy perturbation iteration technique

$$(10) \quad x_{n+1}^\delta = x_n^\delta - \sum_{j=1}^N (I - T_n^* T_n)^{j-1} T_n^* (F(x_n^\delta) - y^\delta).$$

Notably, (10) can alternatively be understood as the N -steps conventional Landweber iteration for resolving the linearized issue [8] as follows: $F(x_n^\delta) + T_n(x - x_n^\delta) = y^\delta$. It is possible to obtain the classical Landweber iteration (4) by using the one-order approximation truncation ($N = 1$). The homotopy perturbation iteration in [13] can be produced using the two-order approximation truncation ($N = 2$):

$$(11) \quad x_{n+1}^\delta = x_n^\delta - (2I - T_n^* T_n) T_n^* (F(x_n^\delta) - y^\delta).$$

It is demonstrated that, in comparison to (4), just half-time is required for (11). It was then successfully used for the inversion of the well log restricted seismic waveform [6]. The convergence study in [14] was conducted with respect to the stopping rule (6) and the nonlinearity condition (5). Additionally, they calculated the convergence rate based on assumption (7), (8) and (9).

It should be emphasised that each repetition step of the procedures in (4) and (11) necessitates the computation of the Fréchet derivative. Under the source condition (7), which is dependent upon the unknown answer x^\dagger , the rate of convergence for each of the aforementioned approaches has been determined. Numerous scholars have explored various versions of the simplified iterative approach and have solved this issue (cf. [10], [15], [16], [17], [18] and [19]). The computation cost of the technique decreases dramatically when compared to the simplified version of the method, which just computes the Fréchet derivative at the initial guess x_0 for the exact solution x^\dagger .

Driven by this benefit, we present a simplified version of the Homotopy perturbation iteration (11), denoted as

$$(12) \quad x_{n+1}^\delta = x_n^\delta - (2I - A_0^* A_0) A_0^* (F(x_n^\delta) - y^\delta),$$

where $A_0 = F'(x_0)$ and A_0^* indicate adjoint of the operator A_0 . While the approach (11) includes the derivative at each iterate step x_n , our method just

uses the Fréchet derivative at the initial guess x_0 . Thus, our approach simplifies the assumptions made in [13] and [14] while simultaneously reducing the computational cost. In this paper, we will examine the convergence analysis of the method (12), using the Morozov type discrepancy principle and appropriate assumptions on the nonlinear operator F . The method's rate of convergence under the Hölder-type source condition will also be examined. Finally, we will provide a numerical example to validate our approach.

2. CONVERGENCE ANALYSIS OF THE METHOD

To prove the method's convergence, we utilise the following assumptions on the nonlinear operator F .

ASSUMPTION 1. (i) For every $x^\dagger \in B_\rho(x_0) \subseteq D(F)$, where x^\dagger is the solution to (1), there exists $\rho > 0$.

(ii) The nonlinear operator is scaled appropriately, i.e.,

$$(13) \quad \|F'(x)\| \leq \frac{1}{2}, \quad x \in B_\rho(x_0)$$

holds.

(iii) The local property

$$(14) \quad \|F(x) - F(x') - F'(x_0)(x - x')\| \leq \eta \|F(x) - F(x')\|,$$

is satisfied by the operator F in a ball $B_\rho(x_0)$, where $\eta < 1$ and $x, x' \in B_\rho(x_0)$.

A number of publications utilise assumptions similar to Assumption 1 (iii) for the convergence analysis of ill-posed equations (cf. [7], [11], [15], [16], [18], [19], [21]). Assumption 1 (iii) can be understood as the tangential cone condition on the operator F .

From equation (14), the triangle inequality yields the following immediately: for every $x, x' \in B_\rho(x_0)$

$$(15) \quad \frac{1}{1+\eta} \|F'(x_0)(x - x')\| \leq \|F(x) - F(x')\| \leq \frac{1}{1-\eta} \|F'(x_0)(x - x')\|.$$

We obtain the inequality from (13)

$$(16) \quad \|I - F'(x_0)^* F'(x_0)\| \leq 1.$$

When dealing with noisy data, the iterations x_n are unable to converge but can nevertheless yield a stable approximation of x^\dagger , as long as they are terminated using the Morozov-type stopping criterion after n_δ steps, i.e.,

$$(17) \quad \|F(x_{n_\delta}^\delta) - y^\delta\| \leq \tau \delta < \|F(x_n^\delta) - y^\delta\|, \quad 0 \leq n < n_\delta$$

where $\tau > 2$ is a positive constant depending on η .

The monotonicity of the iteration error is shown by the following theorem.

THEOREM 1. Let us assume that x^\dagger is a solution of (1) in $B_\rho(x_0)$, and for the situation of the perturbed data y^δ satisfying (2), the iteration is ended

after n_δ steps in accordance with the stopping rule (17), where $\tau > \frac{8(\eta+1)}{3-8\eta}$. If equations (13), (14) and (16) are true, with $0 < \eta < \frac{3}{8}$, we have

$$(18) \quad \|x^\dagger - x_{n+1}^\delta\| \leq \|x^\dagger - x_n^\delta\|, \quad 0 \leq n \leq n_\delta,$$

and if $\delta = 0$,

$$(19) \quad \sum_{n=0}^{\infty} \|F(x_n) - y\|^2 < \infty.$$

Proof. Given $A_0 = F'(x_0)$, $s_n = F(x_n^\delta) - y^\delta$, and $0 \leq n < n_\delta$, we may infer by induction that $x_n^\delta \in B_\rho(x_0)$ from (12), (13), (14) and (16). Consider

$$\begin{aligned} & \|x^\dagger - x_{n+1}^\delta\|^2 - \|x^\dagger - x_n^\delta\|^2 = \\ & = \|x_{n+1}^\delta - x_n^\delta\|^2 + 2\langle x^\dagger - x_n^\delta, x_n^\delta - x_{n+1}^\delta \rangle \\ & = \|(A_0^* A_0 - 2I)A_0^* s_n\|^2 + 2\langle x^\dagger - x_n^\delta, (2I - A_0^* A_0)A_0^*(F(x_n) - y^\delta) \rangle \\ & = \|(A_0^* A_0 - 2I)A_0^* s_n\|^2 + 2\langle F(x_n^\delta) - F(x^\dagger) - A_0(x_n^\delta - x^\dagger) + y - y^\delta, s_n^\delta \rangle \\ & \quad + 2\langle F(x_n^\delta) - F(x^\dagger) - A_0(x_n^\delta - x^\dagger) + y - y^\delta, (I - A_0 A_0^*)s_n^\delta \rangle \\ & \quad - 2\langle s_n^\delta, s_n^\delta \rangle + 2\langle s_n^\delta, (A_0 A_0^* - I)s_n^\delta \rangle \\ & \leq 4\delta(1 + \eta)\|s_n^\delta\| + (4\eta - \frac{3}{2})\|s_n^\delta\|^2. \end{aligned}$$

The right-hand side is negative due to (17) by the stated condition, $0 < \eta < \frac{3}{8}$, $\tau > \frac{8(\eta+1)}{3-8\eta}$, and $n < n_\delta$. We have confirmed (18) and demonstrated the monotonically declining nature of the iteration error. In fact, we have confirmed that the inequality

$$\|x^\dagger - x_{n+1}^\delta\|^2 + (\frac{3}{2} - 4\eta)\|F(x_n) - y\|^2 \leq \|x^\dagger - x_n^\delta\|^2$$

is true for every $n \in \mathcal{N}_0$ if $\delta = 0$.

Consequently, using induction, we obtain

$$\sum_{n=0}^{\infty} \|F(x_n) - y\|^2 \leq \frac{2}{3-8\eta}\|x^\dagger - x_0\|^2.$$

The proof is now complete. \square

REMARK 2. In the event that $\delta \neq 0$, we can demonstrate that

$$(20) \quad n_\delta(\tau\delta)^2 \leq \sum_{n=0}^{n_\delta-1} \|F(x_n^\delta) - y^\delta\|^2 \leq \frac{2\tau}{(3-8\eta)\tau-8(1+\eta)}\|x^\dagger - x_0\|^2.$$

Thus, a well-defined stopping index $n_\delta < \infty$ is determined using the discrepancy principle (17) with $\tau > \frac{8(\eta+1)}{3-8\eta}$.

THEOREM 3. If (1) can be solved in $B_\rho(x_0)$ and Assumption 1 holds, then x_n converges to a solution $x^\dagger \in B_\rho(x_0)$.

Proof. Let x^\dagger be any solution of (1) in $B_\rho(x_0)$, and put

$$(21) \quad r_n = x^\dagger - x_n.$$

$\|r_n\|$ is monotonically decreasing to some $\epsilon \geq 0$ according to Theorem 1. Next, we demonstrate that r_n is a Cauchy sequence. In the case of $i \geq n$, we select m such that $i \geq m \geq n$ and

$$(22) \quad \|F(x_m) - y\| \leq \|F(x_j) - y\|, \quad n \leq j \leq i.$$

Firstly, we have

$$(23) \quad \|r_i - r_n\| \leq \|r_i - r_m\| + \|r_m - r_n\|$$

and

$$(24) \quad \begin{aligned} \|r_i - r_m\|^2 &= 2\langle r_m - r_i, r_m \rangle + \|r_i\|^2 - \|r_m\|^2, \\ \|r_m - r_n\|^2 &= 2\langle r_m - r_n, r_m \rangle + \|r_n\|^2 - \|r_m\|^2. \end{aligned}$$

The final two terms of (24) on both right sides converge to zero for $n \rightarrow \infty$. Now, we use equations (12) and (15) to demonstrate how, as n approaches ∞ , $\langle r_m - r_n, r_m \rangle$ also reduces to zero:

$$\begin{aligned} |\langle r_m - r_n, r_m \rangle| &= \\ &= \left| \sum_{j=n}^{m-1} \langle (2I - A_0^* A_0) A_0^* (F(x_j) - y), r_m \rangle \right| \\ &\leq \sum_{j=n}^{m-1} \left| \langle (2I - A_0 A_0^*) (F(x_j) - y), A_0(x^\dagger - x_m) \rangle \right| \\ &\leq \sum_{j=n}^{m-1} \|(2I - A_0 A_0^*) (F(x_j) - y)\| \|A_0(x^\dagger - x_j + x_j - x_m)\| \\ &\leq \sum_{j=n}^{m-1} \|(2I - A_0 A_0^*) (F(x_j) - y)\| (\|A_0(x^\dagger - x_j)\| + \|A_0(x_j - x_m)\|) \\ &\leq (1 + \eta) \sum_{j=n}^{m-1} \|(2I - A_0 A_0^*) (F(x_j) - y)\| (\|y - F(x_j)\| + \|F(x_m) - F(x_j)\|) \\ &\leq (1 + \eta) \sum_{j=n}^{m-1} \|(2I - A_0 A_0^*) (F(x_j) - y)\| (2\|y - F(x_j)\| + \|F(x_m) - y\|) \\ &\leq 3(1 + \eta) \sum_{j=n}^{m-1} \|(2I - A_0 A_0^*) (F(x_j) - y)\| \|F(x_j) - y\| \\ &\leq 6(1 + \eta) \sum_{j=n}^{m-1} \|(F(x_j) - y)\|^2. \end{aligned}$$

Likewise, it may be demonstrated that

$$\left| \langle r_i - r_m, r_m \rangle \right| \leq 6(1 + \eta) \sum_{j=m}^{i-1} \| (F(x_j) - y) \|^2.$$

From (19) it can be inferred that when $n \rightarrow \infty$, the right-hand side of (24) goes to zero, and from (23), we may deduce that r_n and hence x_n are Cauchy sequences. We designate x^\dagger as the limit of x_n and observe that, as n approaches ∞ , the residuals $F(x_n) - y$ converge to zero, indicating that x^\dagger is a solution of (1). \square

Our subsequent findings indicate that the Landweber iteration becomes a regularization technique as a result of this stopping rule.

THEOREM 4. *Let the assumptions of Theorem 3 hold. If y^δ satisfies (2) and the iteration (12) is stopped in accordance with the stopping rule (17), then the iterates $x_{n_\delta}^\delta$ converge to the solution of (1) as $\delta \rightarrow 0$.*

The proof has a resemblance to [7, Theorem 2.4].

3. CONVERGENCE RATES

We calculate the suggested iteration's rate of convergence in this section. The rate of convergence for the iteration (11) for the source conditions (7) was found in [14]. We observe that, from a practical standpoint, it is quite challenging to validate such assumptions [22]. We use the following source condition to find the convergent rate of the method:

$$(25) \quad x^\dagger - x_0 = (F'(x_0)^* F'(x_0))^\nu w, \quad \nu > 0, \quad w \in X,$$

where $\|w\|$ is small enough.

These approximations will be applied in this section to get the method's convergence rate outcome.

LEMMA 5 (cf. [11]). *Let $A : X \rightarrow Y$ be a bounded linear operator with the property $\|A\| < 1$ and let $\nu \in [0, 1]$. Then*

- (i) $\|(I - A^*A)^n (A^*A)^\nu\| \leq (n+1)^{-\nu},$
- (ii) $\|(I - A^*A)^n A^*\| \leq (n+1)^{-\frac{1}{2}},$
- (iii) $\|\sum_{i=0}^{n-1} (I - A^*A)^i (A^*A)^\nu\| \leq n^{1-\nu}.$

LEMMA 6 (cf. [11]). *Assume p and q are positive. Then, independent of n , there is a positive constant $c(p, q)$ such that*

$$\sum_{i=0}^{n-1} (i+1)^{-p} (n-i)^{-q} \leq c(p, q) (n+1)^{1-p-q} \begin{cases} 1 & \max\{p, q\} < 1 \\ \log(n+1) & \max\{p, q\} = 1 \\ (n+1)^{\max\{p, q\}-1} & \max\{p, q\} > 1. \end{cases}$$

THEOREM 7. *Assume that F satisfies (13) and (14), that y^δ satisfies (2), and that problem (1) has a solution in $B_\rho(x_0)$. A positive constant \mathfrak{C} relying*

exclusively on ν exists if $x^\dagger - x_0$ satisfies (25) with $0 < \nu \leq 1$ and $\|w\|$ is small enough. This can be demonstrated by

$$(26) \quad \|x^\dagger - x_n^\delta\| \leq \mathfrak{C}\|w\|(n+1)^{-\nu},$$

$$(27) \quad \|A_0 e_n^\delta\| \leq \mathfrak{C}\|w\|(n+1)^{-\nu-\frac{1}{2}},$$

for $0 \leq n < n_\delta$. As previously, n_δ is the stopping index of the discrepancy principle (17) in this case, where $\tau > \frac{8(\eta+1)}{3-8\eta}$. For all $n \geq 0$, (26) and (27) hold in the case of exact data ($\delta = 0$).

Proof. The iteration (12) is well-defined according to Theorem 1 since iterates x_n^δ , where $0 \leq n \leq n_\delta$, always remain in $B_\rho(x_0) \subset D(F)$. Furthermore, when $\delta > 0$, the stopping index n_δ is finite. Substitute $e_n^\delta := x^\dagger - x_n^\delta$. Given $0 \leq n < n_\delta$, we get the representation from (12)

$$\begin{aligned} e_{n+1}^\delta &= e_n^\delta - (x_{n+1}^\delta - x_n^\delta) \\ &= (I - 2A_0^*A_0)e_n^\delta - 2A_0^*(y - F(x_n^\delta) - A_0(x^\dagger - x_n^\delta)) \\ &\quad + 2A_0^*(y - y^\delta) + A_0^*A_0A_0^*(y^\delta - F(x_n^\delta)) \\ &= (I - 2A_0^*A_0)e_n^\delta + 2A_0^*z_n^\delta + 2A_0^*(y - y^\delta) + A_0^*A_0A_0^*r_n^\delta, \end{aligned}$$

where $z_n^\delta = -(y - F(x_n^\delta) - A_0(x^\dagger - x_n^\delta))$ and $r_n^\delta = y^\delta - F(x_n^\delta)$.

This produces the closed expression for the error for $0 \leq n < n_\delta$

$$\begin{aligned} e_n^\delta &= (I - 2A_0^*A_0)^n e_0 + \sum_{j=0}^{n-1} (I - 2A_0^*A_0)^j 2A_0^*z_{n-j-1}^\delta + \\ &\quad + \sum_{j=0}^{n-1} (I - 2A_0^*A_0)^j 2A_0^*(y - y^\delta) \\ &\quad + \sum_{j=0}^{n-1} (I - 2A_0^*A_0)^j (A_0^*A_0)A_0^*r_{n-j-1}^\delta, \end{aligned} \quad (28)$$

and consequently

$$\begin{aligned} A_0 e_n^\delta &= (I - 2A_0^*A_0)^n A_0 e_0 + \sum_{j=0}^{n-1} (I - 2A_0^*A_0)^j 2A_0 A_0^*z_{n-j-1}^\delta + \\ &\quad + \sum_{j=0}^{n-1} (I - 2A_0^*A_0)^j 2A_0 A_0^*(y - y^\delta) \\ &\quad + \sum_{j=0}^{n-1} (I - 2A_0^*A_0)^j A_0 (A_0^*A_0)A_0^*r_{n-j-1}^\delta. \end{aligned}$$

We employ induction to demonstrate the result for $0 \leq n < n_\delta$. The proof is simple for $n = 0$, and we take it for granted that the result holds for any j such that $0 \leq j < n$, where $n < n_\delta$.

For $n < n_\delta$,

$$\begin{aligned}
 (29) \quad \|e_n^\delta\| &\leq \|(I - 2A_0^*A_0)^n(A_0^*A_0)^\nu w\| + \sum_{j=0}^{n-1} \|(I - 2A_0^*A_0)^j 2A_0^*\| \|z_{n-j-1}^\delta\| \\
 &\quad + \left\| \sum_{j=0}^{n-1} (I - 2A_0^*A_0)^j 2A_0^*(y - y^\delta) \right\| \\
 &\quad + \sum_{j=0}^{n-1} \|(I - 2A_0^*A_0)^j (A_0^*A_0)A_0^*\| \|r_{n-j-1}^\delta\| \\
 (30) \quad &\leq 2^{-\nu}(n+1)^{-\nu}\|w\| + \sum_{j=0}^{n-1} \sqrt{2}(j+1)^{-1/2} \|z_{n-j-1}^\delta\| + \sqrt{2n}\delta \\
 &\quad + \sum_{j=0}^{n-1} 2^{-3/2}(j+1)^{-3/2} \|r_{n-j-1}^\delta\|.
 \end{aligned}$$

and

$$\begin{aligned}
 \|A_0 e_n^\delta\| &\leq \|(I - 2A_0^*A_0)^n A_0 e_0\| + \sum_{j=0}^{n-1} \|(I - 2A_0^*A_0)^j 2A_0 A_0^*\| \|z_{n-j-1}^\delta\| \\
 &\quad + \left\| \sum_{j=0}^{n-1} (I - 2A_0^*A_0)^j 2A_0 A_0^*(y - y^\delta) \right\| \\
 &\quad + \sum_{j=0}^{n-1} \|(I - 2A_0^*A_0)^j A_0 (A_0^*A_0)A_0^*\| \|r_{n-j-1}^\delta\| \\
 (31) \quad &\leq 2^{-\nu-\frac{1}{2}}(n+1)^{-\nu-\frac{1}{2}}\|w\| \\
 &\quad + \sum_{j=0}^{n-1} (j+1)^{-1} \|z_{n-j-1}^\delta\| + \delta + 2^{-2} \sum_{j=0}^{n-1} (j+1)^{-2} \|r_{n-j-1}^\delta\|
 \end{aligned}$$

Using the triangle inequality, equations (15), (17) and the induction assumption, we now obtain

$$(32) \quad \|y^\delta - F(x_n^\delta)\| \leq 2\|y - F(x_n^\delta)\| \leq \frac{2}{1-\eta} \|A_0(x^\dagger - x_n^\delta)\| \leq \frac{2}{1-\eta} \mathfrak{C}\|w\|(n+1)^{-\nu-\frac{1}{2}},$$

and

$$\begin{aligned}
 \|z_n^\delta\| &\leq \eta\|y - F(x_n^\delta)\| \\
 &\leq \frac{\eta}{1-\eta} \|A_0 e_n^\delta\| \\
 (33) \quad &\leq \frac{\eta}{1-\eta} \mathfrak{C}\|w\|(n+1)^{-\nu-\frac{1}{2}}
 \end{aligned}$$

Consequently,

$$\sum_{j=0}^{n-1} \sqrt{2}(j+1)^{-\frac{1}{2}} \|z_{n-j-1}^\delta\| \leq \sqrt{2} \frac{\eta}{1-\eta} \mathfrak{C} \|w\| \sum_{j=0}^{n-1} (j+1)^{-\frac{1}{2}} (n-j)^{-\nu-\frac{1}{2}}$$

and

$$\sum_{j=0}^{n-1} 2^{\frac{-3}{2}} (j+1)^{-\frac{3}{2}} \|r_{n-j-1}^\delta\| \leq \frac{2^{\frac{-1}{2}}}{1-\eta} \mathfrak{C} \|w\| \sum_{j=0}^{n-1} (j+1)^{-\frac{3}{2}} (n-j)^{-\nu-\frac{1}{2}}$$

Thus, applying [Lemma 6](#), we obtain

$$(34) \quad \sum_{j=0}^{n-1} \sqrt{2}(j+1)^{-\frac{1}{2}} \|z_{n-j-1}^\delta\| \leq a_\nu \|w\| (n+1)^{-\nu}$$

and

$$(35) \quad \sum_{j=0}^{n-1} 2^{\frac{-3}{2}} (j+1)^{-\frac{3}{2}} \|r_{n-j-1}^\delta\| \leq b_\nu \|w\| (n+1)^{-\nu}$$

where the constants $a_\nu > 0$ and $b_\nu > 0$ are dependent on ν . Therefore

$$\begin{aligned} \|e_n^\delta\| &\leq 2^{-\nu} (n+1)^{-\nu} \|w\| + a_\nu \|w\| (n+1)^{-\nu} + \sqrt{2n\delta} + b_\nu \|w\| (n+1)^{-\nu} \\ &\leq (2^{-\nu} + a_\nu + b_\nu) \|w\| (n+1)^{-\nu} + \sqrt{2n\delta}. \end{aligned}$$

Likewise, one may demonstrate that

$$\begin{aligned} \|A_0 e_n^\delta\| &\leq \\ &\leq 2^{-\nu-\frac{1}{2}} (n+1)^{-\nu-\frac{1}{2}} \|w\| + \tilde{a}_\nu \|w\| (n+1)^{-\nu-\frac{1}{2}} + \delta + \tilde{b}_\nu \|w\| (n+1)^{-\nu-\frac{1}{2}} \\ (36) \quad &\leq (2^{-\nu-\frac{1}{2}} + \tilde{a}_\nu + \tilde{b}_\nu) \|w\| (n+1)^{-\nu-\frac{1}{2}} + \delta. \end{aligned}$$

Here, $\tilde{a}_\nu > 0$ and $\tilde{b}_\nu > 0$ depends on ν .

Owing to (17), $\tau > \frac{8(\eta+1)}{3-8\eta}$, as follows:

$$\frac{8(\eta+1)}{3-8\eta} \delta \leq \tau \delta \leq \|F(x_n^\delta) - y^\delta\| \leq \frac{1}{1-\eta} \|A_0 e_n^\delta\| + \delta.$$

Therefore, using the above result (36), we obtain

$$\frac{8(\eta+1)}{3-8\eta} \delta \leq \frac{1}{1-\eta} (2^{-\nu-\frac{1}{2}} + \tilde{a}_\nu + \tilde{b}_\nu) \|w\| (n+1)^{-\nu-\frac{1}{2}} + \frac{2-\eta}{1-\eta} \delta.$$

This would provide

$$(37) \quad \delta \leq \frac{(3-8\eta)}{-16\eta^2+19\eta+2} (2^{-\nu-\frac{1}{2}} + \tilde{a}_\nu + \tilde{b}_\nu) \|w\| (n+1)^{-\nu-\frac{1}{2}}.$$

Therefore,

$$(38) \quad \|e_n^\delta\| \leq \mathfrak{C} \|w\| (n+1)^{-\nu},$$

and

$$(39) \quad \|A_0 e_n^\delta\| \leq \mathfrak{C} \|w\| (n+1)^{-\nu-\frac{1}{2}},$$

where

$$\mathfrak{C} = \max \left(\left(1 + \frac{(3-8\eta)}{-16\eta^2+19\eta+2} \right) (2^{-\nu-\frac{1}{2}} + \tilde{a}_\nu + \tilde{b}_\nu), \right. \\ \left. 2^{-\nu} + a_\nu + b_\nu + \sqrt{2} \left(\frac{(3-8\eta)}{-16\eta^2+19\eta+2} \right) (2^{-\nu-\frac{1}{2}} + \tilde{a}_\nu + \tilde{b}_\nu) \right).$$

□

THEOREM 8. Assuming the conditions of [Theorem 7](#), we obtain

$$(40) \quad n_\delta \leq \mathfrak{C}_1 \left(\frac{\|w\|}{\delta} \right)^{\frac{2}{2\nu+1}},$$

and

$$(41) \quad \|x^\dagger - x_{n_\delta}^\delta\| \leq \mathfrak{C}_2 \|w\|^{\frac{1}{2\nu+1}} \delta^{\frac{2\nu}{2\nu+1}},$$

where \mathfrak{C}_1 and $\mathfrak{C}_2 > 0$ are positive constant that depends exclusively on ν .

Proof. Equation (28) allows us to write

$$\begin{aligned} e_{n_\delta}^\delta &= (I - 2A_0^*A_0)^{n_\delta} e_0 + \\ &+ \sum_{j=0}^{n_\delta-1} (I - 2A_0^*A_0)^j 2A_0^* z_{n_\delta-j-1}^\delta + \sum_{j=0}^{n_\delta-1} (I - 2A_0^*A_0)^j 2A_0^* (y - y^\delta) \\ &+ \sum_{j=0}^{n_\delta-1} (I - 2A_0^*A_0)^j (A_0^*A_0) A_0^* r_{n_\delta-j-1}^\delta, \\ &= (A_0^*A_0)^\nu W_{n_\delta} + \sum_{j=0}^{n_\delta-1} (I - 2A_0^*A_0)^j 2A_0^* (y - y^\delta), \end{aligned}$$

where $W_{n_\delta} = (I - 2A_0^*A_0)^{n_\delta} w + \sum_{j=0}^{n_\delta-1} 2^{\frac{1}{2}+\nu} (I - 2A_0^*A_0)^j (2A_0^*A_0)^{\frac{1}{2}-\nu} \tilde{z}_{n_\delta-j-1}^\delta + \sum_{j=0}^{n_\delta-1} 2^{\frac{-3}{2}+\nu} (I - 2A_0^*A_0)^j (2A_0^*A_0)^{\frac{3}{2}-\nu} \tilde{r}_{n_\delta-j-1}^\delta$, we can write $\|\tilde{z}_j^\delta\| = \|z_j^\delta\|$ and $\|\tilde{r}_j^\delta\| = \|r_j^\delta\|$, $j = 0, 1, \dots, n_\delta - 1$.

$$\begin{aligned} \|W_{n_\delta}\| &\leq \|(I - 2A_0^*A_0)^{n_\delta} w\| + \sum_{j=0}^{n_\delta-1} \|2^{\frac{1}{2}+\nu} (I - 2A_0^*A_0)^j (2A_0^*A_0)^{\frac{1}{2}-\nu}\| \|z_{n_\delta-j-1}^\delta\| \\ &+ \sum_{j=0}^{n_\delta-1} \|2^{\frac{-3}{2}+\nu} (I - 2A_0^*A_0)^j (2A_0^*A_0)^{\frac{3}{2}-\nu}\| \|r_{n_\delta-j-1}^\delta\| \\ &\leq (n_\delta + 1)^0 \|w\| + \sum_{j=0}^{n_\delta-1} 2^{\frac{1}{2}+\nu} (j+1)^{\nu-\frac{1}{2}} \|z_{n_\delta-j-1}^\delta\| \\ &+ \sum_{j=0}^{n_\delta-1} 2^{\frac{-3}{2}+\nu} (j+1)^{\nu-\frac{3}{2}} \|r_{n_\delta-j-1}^\delta\| \end{aligned}$$

We have

$$\sum_{j=0}^{n_\delta-1} 2^{\frac{1}{2}+\nu}(j+1)^{\nu-\frac{1}{2}} \|z_{n_\delta-j-1}^\delta\| \leq \frac{2^{\frac{1}{2}+\nu}\eta}{1-\eta} \mathfrak{C} \|w\| \sum_{j=0}^{n_\delta-1} (j+1)^{\nu-\frac{1}{2}} (n_\delta-j)^{-\nu-\frac{1}{2}}$$

and

$$\sum_{j=0}^{n_\delta-1} 2^{\frac{-3}{2}+\nu}(j+1)^{\nu-\frac{3}{2}} \|r_{n_\delta-j-1}^\delta\| \leq \frac{2^{\frac{-1}{2}+\nu}}{1-\eta} \mathfrak{C} \|w\| \sum_{j=0}^{n_\delta-1} (j+1)^{\nu-\frac{3}{2}} (n_\delta-j)^{-\nu-\frac{1}{2}}$$

By applying [Lemma 6](#),

$$\sum_{j=0}^{n_\delta-1} 2^{\frac{1}{2}+\nu}(j+1)^{\nu-\frac{1}{2}} \|z_{n_\delta-j-1}^\delta\| \leq c_\nu (n_\delta+1)^0 \|w\|$$

and

$$\sum_{j=0}^{n_\delta-1} 2^{\frac{-3}{2}+\nu}(j+1)^{\nu-\frac{3}{2}} \|r_{n_\delta-j-1}^\delta\| \leq \tilde{c}_\nu (n_\delta+1)^{-1} \|w\|$$

Since we are aware that $\|w\|$ must be small, we take $\|w\| \leq 1$. Thus, we have

$$\begin{aligned} \|W_{n_\delta}\| &\leq (n_\delta+1)^0 \|w\| + c_\nu (n_\delta+1)^0 \|w\| + \tilde{c}_\nu (n_\delta+1)^{-1} \|w\| \\ &\leq (1 + c_\nu + \tilde{c}_\nu) \|w\|. \end{aligned}$$

Therefore,

$$\begin{aligned} \|A_0(A_0^*A_0)^\nu W_{n_\delta}\| &= \|A_0 e_{n_\delta}^\delta - (I - (I - A_0 A_0^*)^k)(y^\delta - y)\| \\ &\leq \|A_0 e_{n_\delta}^\delta\| + \delta \\ &\leq \|F(x_{n_\delta}) - F(x^\dagger) - A_0 e_{n_\delta}^\delta\| + \|F(x_{n_\delta}^\delta) - F(x^\dagger)\| + \delta \\ &\leq (\eta+1) \|y - F(x_{n_\delta}^\delta)\| + \delta \\ &\leq ((\eta+1)(\tau+1)+1)\delta. \end{aligned}$$

Applying the inequality of interpolation, we obtain

$$\begin{aligned} \|(A_0^*A_0)^\nu W_{n_\delta}\| &\leq \left(((\eta+1)(\tau+1)+1)\delta \right)^{\frac{2\nu}{2\nu+1}} \left((1 + c_\nu + \tilde{c}_\nu) \|w\| \right)^{\frac{1}{2\nu+1}} \\ &\leq \mathfrak{D} \delta^{\frac{2\nu}{2\nu+1}} \|w\|^{\frac{1}{2\nu+1}}, \end{aligned}$$

where \mathfrak{D} is positive constant. For $n_\delta = 0$,

$$(42) \quad \|e_{n_\delta}^\delta\| \leq \mathfrak{D} \|w\|^{\frac{1}{2\nu+1}} \delta^{\frac{2\nu}{2\nu+1}},$$

and when $n_\delta > 0$, we apply [\(37\)](#) with $n = n_\delta - 1$ to obtain

$$(43) \quad \delta \leq \Gamma \|w\| n_\delta^{-\nu-\frac{1}{2}},$$

where $\Gamma = \frac{(3-8\eta)}{-16\eta^2+19\eta+2} (2^{-\nu-\frac{1}{2}} + \tilde{a}_\nu + \tilde{b}_\nu)$, and hence

$$n_\delta \leq \mathfrak{C}_1 \left(\frac{\|w\|}{\delta} \right)^{\frac{2}{2\nu+1}},$$

where $\mathfrak{C}_1 = \Gamma^{\frac{2}{2\nu+1}}$.

Using the outcome that we obtain,

$$\begin{aligned} \|e_{n_\delta}^\delta\| &\leq \|(A_0^* A_0)^\nu W_{n_\delta}\| + \sqrt{2n_{n_\delta}} \delta \\ &\leq \mathfrak{D} \|w\|^{\frac{1}{2\nu+1}} \delta^{\frac{2\nu}{2\nu+1}} + \sqrt{2\mathfrak{C}_1} \|w\|^{\frac{1}{2\nu+1}} \delta^{1-\frac{1}{2\nu+1}} \\ &\leq \mathfrak{C}_2 \|w\|^{\frac{1}{2\nu+1}} \delta^{\frac{2\nu}{2\nu+1}}, \end{aligned}$$

where $\mathfrak{C}_2 = \mathfrak{D} + \sqrt{2\mathfrak{C}_1}$. \square

REMARK 9. *In contrast to the suggested simplified HPI (12), the HPI (11) taken into consideration in [14] necessitates the extra condition (8) and (9) on the non-linear operator F in order to show the convergence rate conclusion. The convergence rate result cannot be maintained in practical problems if the operator does not meet condition (8) and (9). Here, we give an example that does not meet assumptions (8) and (9).*

4. NUMERICAL EXAMPLE

This section examines a numerical example to demonstrate the adaptability of the suggested simplified homotopy method. Matlab is used for numerical calculations.

Here, we look at the nonlinear model problem, which is recovering the parameter estimation problem's diffusion term. Let $f \in L^2(\Omega)$ and $\Omega \subset \mathbb{R}^d (d = 1, 2)$ represent an open bounded domain with a Lipschitz border Γ . We examine calculating the diffusion coefficient c in equation

$$(44) \quad \begin{cases} -(c(t)u(t)_t)_t = f(t), & \text{in } \Omega \\ u = 0, & \text{on } \Gamma. \end{cases}$$

We consider that $L^2(\Omega)$ contains the exact diffusion coefficient c^\dagger . In the domain $D(F) := \{c \in H^1(\Omega) : c(x) \geq \underline{c} > 0\}$, there is a solution $u = u(c)$ in $H^1(\Omega)$ for every c in the domain. We may define the nonlinear operator $F : X = L^2(\Omega) \rightarrow Y = L^2(\Omega)$ with $F(c) = u(c)$ by using the Sobolev embedding $H^1(\Omega) \rightarrow L^2(\Omega)$. It was demonstrated in [5], [7] and [20] that this operator is Fréchet differentiable with

$$(45) \quad \begin{aligned} F'(c)h &= T_c^{-1}[(hu_t(c))_t], F'(c)^*w = -B^{-1}[u_t(c)(T_c^{-1}(w))_t], \\ c &\in D(F), h, w \in L^2(\Omega), \end{aligned}$$

where $T_c : H^2(\Omega) \cap H_0^1(\Omega) \rightarrow L^2(\Omega)$ is defined by $T_c u := -(c(t)u_t(t))_t$ and $B : D(B) := \{\Psi \in H^2(\Omega) : \Psi' = 0 \text{ on } \Gamma\} \rightarrow L^2(\Omega)$ is defined as $B\Psi := -\Psi'' + \Psi$; note that B^{-1} is the adjoint of the embedding operator from $H^1 \rightarrow L^2$.

For every c in $B_\rho(c_0)$, if $u_t(c) \geq b$, $b > 0$, then (5) holds locally according to Lemma 2.6 in [23], which guarantees the convergence of the HPI (11). Assumption (14) is a particular instance of (5), which guarantees the convergence of the proposed iteration (12) according to the Theorem 3 and Theorem 4.

Convergence rate results do not hold for the HPI (11) since F , regrettably, does not satisfy assumption (8) and (9) (see [10], [23]). We do not need the assumption (8) for the convergence rate results in the suggested simplified HPI (12).

From the iteration (12), for all n , $x_{n+1}^\delta - x_n^\delta \in \mathcal{R}(A_0^*)$. Therefore, in particular $x_{n_\delta}^\delta - x_0 \in \mathcal{R}(A_0^*) \forall \delta$, which means $\lim_{\delta \rightarrow 0} x_{n_\delta}^\delta - x_0 \in \mathcal{R}(A_0^*)$. Hence

$$(46) \quad x^\dagger - x_0 \in \mathcal{R}(A_0^*)$$

Therefore, assumptions (25) is satisfied due to both assumption (14) and (46)

Assume that $\Omega = [0, 1]$. The right-hand side of differential equation (44) is explicitly determined using exact information $u^\dagger(t)$. By introducing random noise to the precise data $u^\dagger(t)$ at a specified noise level δ , the perturbed data $u^\delta(t)$ satisfying $\|u^\delta(t) - u^\dagger(t)\| \leq \delta$ is obtained. On a uniform grid with different grid points (N), the differential equations involving the Fréchet differentiable (45) are solved using the finite element method using linear splines. With $\tau = 5$, the iteration is terminated using the stopping criteria (17).

We use the function

$$f(t) = -e^t(1 + \frac{1}{2} \sin(2\pi t) + \pi \cos(2\pi t)) + (e - 1)\pi \cos(2\pi t),$$

and the exact data $u^\dagger(t) = e^t + (1 - e)t - 1$, so the exact solution is $c^\dagger(t) = 1 + \frac{1}{2} \sin(2\pi t)$.

We start the iteration with the initial guess

$$c_0(t) = 1 + \frac{1}{2} \sin(2\pi t) + 200t^2(1 - t)^2(0.25 - t)(0.75 - t).$$

Then, according to [21]: $c_0 - c^\dagger = 200t^2(1 - t)^2(0.25 - t)(0.75 - t) \in R(F'(x^\dagger)^*)$.

We employ both the suggested simplified Homotopy perturbation iteration (12) and the Homotopy perturbation iteration (11). Several values of δ and grid point N are chosen in order to show how the convergence rates depend on the noise level. The numerical results are given in Table 1 and Table 2. Fig. 1–Fig. 6 respectively, represent graphical results for the methods (11) and (12).

Obviously, Fig. 1–Fig. 6 show that the approximate effect of simplified Homotopy perturbation iteration is better. Compared with Homotopy perturbation iteration, we discover in Table 1 and Table 2 that the error norm of simplified Homotopy perturbation iteration is much smaller within less iteration steps and less computational time, and also the convergent rate of simplified Homotopy perturbation iteration (12) is faster than Homotopy perturbation iteration (11).

5. CONCLUSION

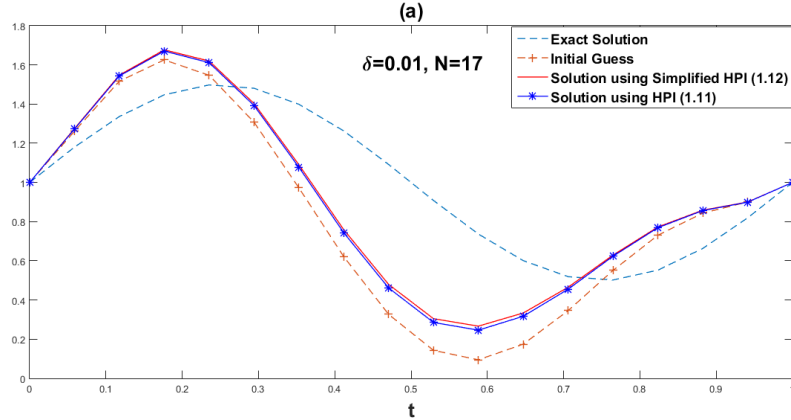
In this study, we have performed the analysis and evaluated a simplified version of the Homotopy perturbation iteration (11). The suggested technique has the advantage of just computing the Fréchet derivative once, rather than

δ	N	n_δ	$Error = \ c_{n_\delta}^\delta - c^\dagger\ $	Time(s)
0.01	17	424	1.3377	1.5375
0.005	17	4853	0.9980	16.9290
0.001	17	53252	0.2844	189.9022
0.0005	17	134204	0.1877	421.6156
0.01	33	1972	1.6762	8.7951
0.005	33	8511	1.1900	31.0274
0.01	65	4457	2.0520	21.7195
0.005	65	13948	1.3607	75.8592

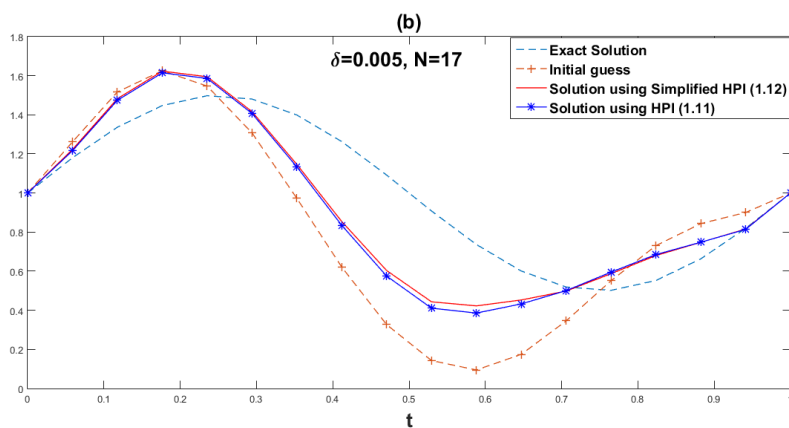
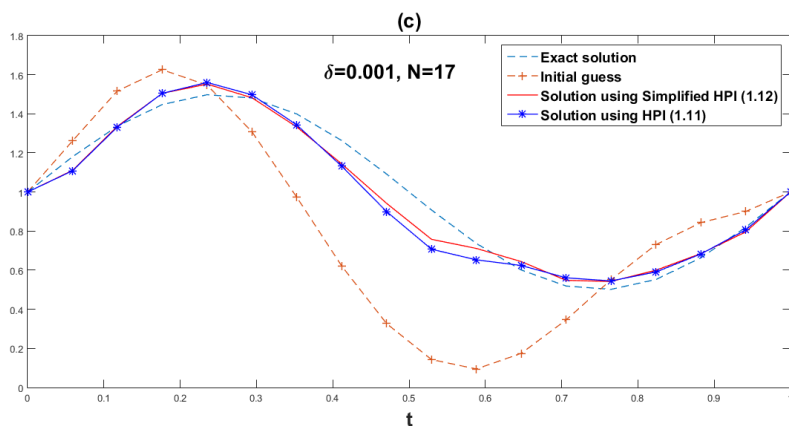
Table 1. The results of the Simplified HPI (12).

δ	N	n_δ	$Error = \ c_{n_\delta}^\delta - c^\dagger\ $	Time(s)
0.01	17	510	1.3569	1.8747
0.005	17	5970	1.0488	18.4612
0.001	17	83441	0.3475	265.8005
0.0005	17	189874	0.2529	628.5853
0.01	33	2350	1.7366	9.1521
0.005	33	10628	1.2667	42.0180
0.01	65	5312	2.1569	26.6291
0.005	65	18020	1.4721	91.1350

Table 2. The results of the HPI (11).


Fig. 1. Solution when $N = 17$ and $\delta = 0.01$

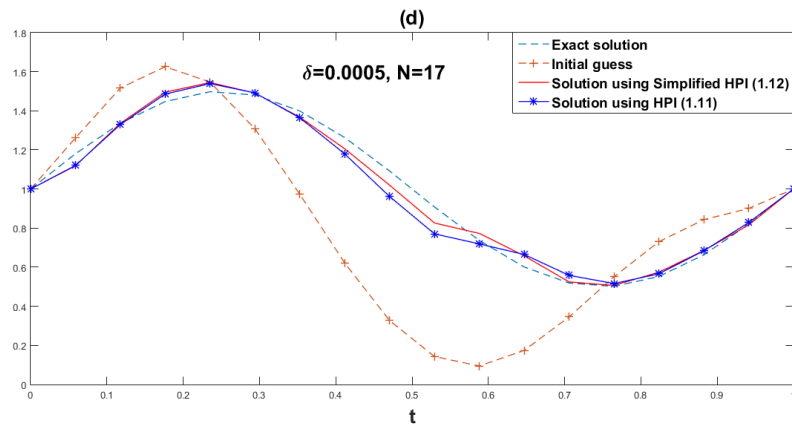
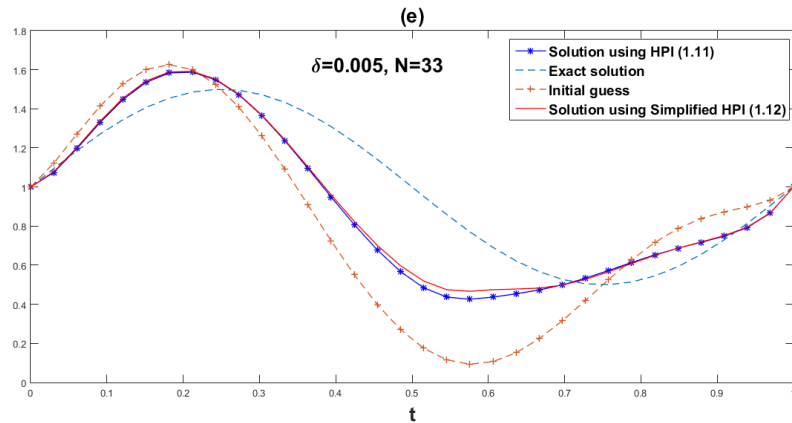
at each iteration step. The calculation in this approach becomes simpler than the traditional Homotopy perturbation iterations, as the iteration (12) and the source condition (25) only entail the Fréchet derivative at initial guess x_0 to the exact solution x^\dagger of (1). The suggested iteration is competitive in terms of reducing the overall computing time as well as error when compared




Fig. 2. Solution when $N = 17$ and $\delta = 0.005$ Fig. 3. Solution when $N = 17$ and $\delta = 0.001$

to the traditional Homotopy perturbation iteration as demonstrated by the numerical example.

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Fig. 4. Solution when $N = 17$ and $\delta = 0.0005$ Fig. 5. Solution when $N = 33$ and $\delta = 0.005$

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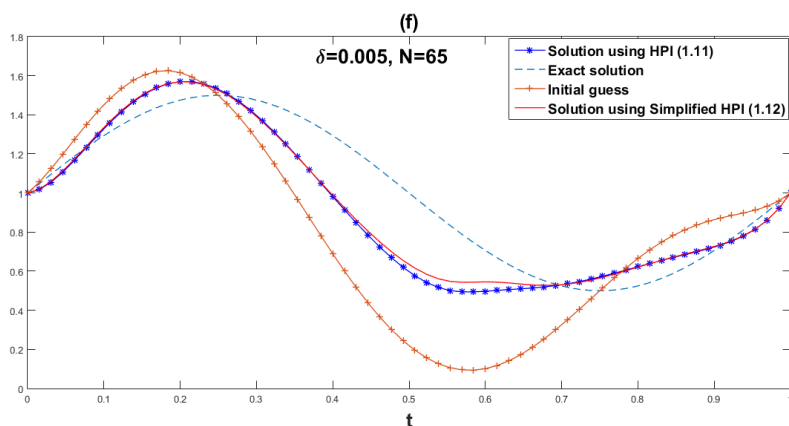



Fig. 6. Solution when $N = 65$ and $\delta = 0.005$

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